Kinematics and Internal Dynamics Formulations of the Canis Compressed Air Engine

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Abstract

Prior to this paper some study was conducted that resulted to the conceptual design of the Canis Compressed Air Engine. This paper brings out the mathematical design of the Canis Compressed Air engine. It also highlights the study of the kinematic and internal dynamic formula. With the engine components movement likened to an eccentric circular cam, with a radial movement (R_i) and a normal movement component (D_i) and a different (d_i), with the magnitude of the radial movement (R_i) and different (d_i) considered constant, respectively 4units and 1units, the mathematical formula describing the displacement, velocity and acceleration were verified. This verification, done graphically, ended up with the conclusion that this design of a compressed air engine, complies theoretically with the recommendations of accurate intake and exhaust positioning, smooth running and high efficiency due to absence of backlashing forces during intake. The theoretical internal dynamics analysis further proofs the susceptibility of a perfect intake then explosive expansion and a perfect exhaust for every expansion chamber.

Keywords: Kinematics, dynamics, eccentric, circular, cam and free-body

1. Introduction

The Canis Compressed Air engine is an engine design that is powered solely by compressed air. This engine uses the rotary engine principle and prior to this paper the engine's physical appearance was designed with the geometry of the parts and their assembly established by this same author (Fru, 2015). This paper analyses the kinematics and dynamics formulations which were graphically studied and conclusions derived on the functioning principles of the machine.

2. Kinematics Formulations of the Canis Compressed – Air (Cc) Engine

In order to realize the kinematics formulations of this mechanism, (Vinogradov, 2000) a skeleton representation of what was thought of as the mechanism was established.

2.1 Skeleton Display of CC Engine

The diagram below represents what was considered as an instant in the movement of the CC engine. There too, has been associated the general axes (x, y) and the local axes (x_1, y_1) of the assembly. (Kovacevic, 2009).

The graphical display brought in some satisfaction on the balancing of the mechanism, for vividly can be seen an annulment of the radial forces and even the radial movements. A complete functional cycle of this design consists of an intake, power (expansion) and exhaust. These were considered to occur for each complete wing cycle of all the expansion chambers and so, to result each to a complete cycle of the output rotor centered at the point 'O'.

It was noted that; EC_5 (Expansion Chamber 5) was at the verge of completing its exhaust thus its wing had little or no influence on the movement of the cylinder, instead, the radial force from EC_2 was compelled to be compensated by the sum of contributions donated by the other normal forces from the different ECs that were not actually in line with their respective radial reaction forces. The annulment of radial forces in this manner insinuates equilibrium in the engine design.



Figure 1. Skeleton display of CC engine

2.2 Kinematic Diagram

The movement of contact points of each wing with the circle (C') was taken for a vector triangle consisting of a radial movement (R_i) and a normal movement component (D_i) and a different (d_i).

From prior knowledge on the CC engine and the above skeleton, the CC engine rotor cylinder was seen to experience an eccentric circular cam movement centered at 'O' and tangent to the circle 'C' all displayed above. This implies that the contact point of each wing with this cam experiences movements which were likened to the loop vector below.



Figure 2. The links of CC engine

The above vector display was then interpreted as follows;

- $D_1 = D_1(t)$ that is, the magnitude of D_1 varies with time but was considered as unknown
- $d_1 = \text{const}$ that is, the magnitude of d_1 is constant with time
- $R_1 = \text{const}$ that is, the magnitude of R_1 is constant with time
- $\theta_1 = \theta_1(t)$ that is, the angle θ_1 varies with time
- $\gamma_1 = \gamma_1(t)$ that is, the angle γ_1 varies with time but was considered as known.
- It was equally considered that the vector D_1 remains normal to our local x_1 -axis and equals the sum of the vectors d_1 and R_1 .

2.3 Displacement Calculations

With these interpretations, the loop closure equation was established to be equivalent to

Equation 1:

$$R_1(\cos\theta_1,\sin\theta_1) + d_1(\cos\gamma_1,\sin\gamma_1) + D_1(\cos 3\pi/2,\sin 3\pi/2) = 0$$

This is an equation with two unknowns D_1 and θ_1 which was solved as follows;

$$\Rightarrow R_1(\cos \theta_1, \sin \theta_1) + D_1(\cos 3\pi/2, \sin 3\pi/2) = -d_1(\cos \gamma_1, \sin \gamma_1)$$

Suppose $(\cos 3\pi/2, \sin 3\pi/2)$ and $(-\sin 3\pi/2, \cos 3\pi/2)$ are unit vectors parallel and perpendicular respectively to the vector D₁. A scalar product, all through, with each of these unit vectors was carried out on (1). For $(\cos 3\pi/2, \sin 3\pi/2)$

$$\Leftrightarrow \mathsf{R}_1(\cos\theta_1,\sin\theta_1).(\cos 3\pi/2,\sin 3\pi/2) + \mathsf{D}_1 = -\mathsf{d}_1(\cos\gamma_1,\sin\gamma_1).(\cos 3\pi/2,\sin 3\pi/2)$$

$$\Rightarrow D_1 + R_1(\cos\theta_1\cos 3\pi/2 + \sin\theta_1\sin 3\pi/2) = -d_1(\cos\gamma_1\cos 3\pi/2 + \sin\gamma_1\sin 3\pi/2)$$

Equation 2:

$$\Rightarrow$$
 D₁ + R₁ cos($\theta_1 - 3\pi/2$) = -d₁ cos($\gamma_1 - 3\pi/2$)

And for $(-\sin 3\pi/2, \cos 3\pi/2)$

$$\Leftrightarrow \mathbf{R}_1(\cos\theta_1,\sin\theta_1).(-\sin 3\pi/2,\cos 3\pi/2) = -\mathbf{d}_1(\cos\gamma_1,\sin\gamma_1).(-\sin 3\pi/2,\cos 3\pi/2)$$
$$\Rightarrow \mathbf{R}_1(\sin\theta_1\cos 3\pi/2 - \cos\theta_1\sin 3\pi/2) = -\mathbf{d}_1(\sin\gamma_1\cos 3\pi/2 - \cos\gamma_1\sin 3\pi/2)$$

Equation 3:

$$\Rightarrow \mathbf{R}_1 \sin(\theta_1 - 3\pi/2) = -\mathbf{d}_1 \sin(\gamma_1 - 3\pi/2)$$

From equation (2) and (3), was deduced

$$R_1 \cos(\theta_1 - 3\pi/2) = -d_1 \cos(\gamma_1 - 3\pi/2) - D_1$$

And

$$R_1 \sin(\theta_1 - 3\pi/2) = -d_1 \sin(\gamma_1 - 3\pi/2)$$

Thus

$$R_1^2 = (d_1 \cos(\gamma_1 - 3\pi/2) + D_1)^2 + (d_1 \sin(\gamma_1 - 3\pi/2))^2$$

$$\Rightarrow D_1^2 + 2D_1 d_1 \cos(\gamma_1 - 3\pi/2) + d_1^2 - R_1^2 = 0$$

Equation 4:

$$\therefore D_1 = -d_1 \cos(\gamma_1 - 3\pi/2) \pm \sqrt{d_1^2 \cos^2(\gamma_1 - 3\pi/2) + R_1^2 - d_1^2}$$

Given that D_1 is always positive therefore;

$$D_{1} = -d_{1}\cos(\gamma_{1} - 3\pi/2) + \sqrt{d_{1}^{2}\cos^{2}(\gamma_{1} - 3\pi/2) + R_{1}^{2} - d_{1}^{2}}$$
$$\therefore D_{1} = d_{1}\sin\gamma_{1} + \sqrt{R_{1}^{2} - d_{1}^{2}\cos^{2}\gamma_{1}}$$

Equation 5:

$$\theta_{1}^{*} = 3\pi/2 + \sin^{-1} \left(\frac{-d_{1} \sin(\gamma_{1} - 3\pi/2)}{R_{1}} \right)$$
$$\therefore \theta_{1}^{*} = 3\pi/2 + \sin^{-1} \left(\frac{-d_{1} \cos\gamma_{1}}{R_{1}} \right)$$
$$\theta_{1} = \begin{cases} \theta_{1}^{*} if - \sin\theta_{1} \ge 0 \text{ and } \cos\theta_{1} < 0 \\ \pi - \theta_{1}^{*} if - \sin\theta_{1} < 0 \text{ and } \cos\theta_{1} < 0 \\ \pi + \theta_{1}^{*} if - \sin\theta_{1} < 0 \text{ and } \cos\theta_{1} \ge 0 \\ 2\pi - \theta_{1}^{*} if - \sin\theta_{1} \ge 0 \text{ and } \cos\theta_{1} \ge 0 \end{cases}$$

Given that actual dimensions of d_1 and R_1 were not yet found, it was supposed that, their respective values were 1 and 4 units respectively, then graphically; the values of D_1 for varying γ_1 and that of θ_1 in radians for varying γ_1 were seen to follow the patterns below.



Figure 3. Displacement diagram of CC engine

The displacement of D_1 is sinusoidal, describing a complete circle in a 360° rotation of γ_1 . The dwelling of θ_1 , during the transitions from the maximum displacement to minimum displacement describing the quasi instantaneous rise or fall, implies that time is given for the gas to expand and the wings to return to their initial position.

Further study on the velocity of the core moving components of the engine was done.

2.4 Velocity Calculations

From the loop closure equation above, the following velocity equation was deduced;

Equation 6:

$$\mathbf{R}_1\dot{\mathbf{\theta}}_1(-\sin\theta_1,\cos\theta_1) + \mathbf{d}_1\dot{\mathbf{\gamma}}_1(-\sin\gamma_1,\cos\gamma_1) + \dot{\mathbf{D}}_1(\cos 3\pi/2,\sin 3\pi/2) = \mathbf{0}$$

It was assumed that $\dot{\gamma_1}$ was known, for it is angular velocity expected from the rotor and that, \dot{D}_1 and $\dot{\theta}_1$ were to be calculated. Using the unit vectors that were given above, scalar products across equation 6 were done:

For $(\cos 3\pi/2, \sin 3\pi/2)$

 $R_1\dot{\theta}_1(\cos\theta_1\sin 3\pi/2 - \sin\theta_1\cos 3\pi/2) = -d_1\dot{\gamma}_1(\cos\gamma_1\sin 3\pi/2 - \sin\gamma_1\cos 3\pi/2) - \dot{D}_1$ Equation 7:

$$\Rightarrow \mathbf{R}_1 \dot{\mathbf{\theta}}_1 \cos \mathbf{\theta}_1 = \dot{\mathbf{D}}_1 - \mathbf{d}_1 \dot{\gamma_1} \cos \gamma_1$$

And for $(-\sin 3\pi/2, \cos 3\pi/2)$

 $R_1\dot{\theta}_1(\sin\theta_1\sin3\pi/2 + \cos\theta_1\cos3\pi/2) = -d_1\dot{\gamma}_1(\sin\gamma_1\sin3\pi/2 + \cos\gamma_1\cos3\pi/2)$ Equation 8:

$$\Rightarrow$$
 R₁ θ_1 sin $\theta_1 = -d_1\dot{\gamma_1}$ sin γ_1

From equations 7 and 8, was deduced Equation 9:

$$\dot{\theta}_1 = \frac{-d_1 \dot{\gamma}_1 \sin \gamma_1}{R_1 \sin \theta_1}$$

And

$$\dot{D}_1 = R_1 \dot{\theta}_1 \cos \theta_1 + d_1 \dot{\gamma}_1 \cos \gamma_1$$

$$\Rightarrow \dot{D}_1 = \frac{R_1 d_1 \dot{\gamma}_1 \sin \theta_1 \cos \gamma_1 - R_1 d_1 \dot{\gamma}_1 \sin \gamma_1 \cos \theta_1}{R_1 \sin \theta_1}$$

$$\Rightarrow \dot{D}_1 = \frac{d_1 \dot{\gamma}_1 \sin(\theta_1 - \gamma_1)}{\sin \theta_1}$$

Equation 10:

$$\dot{\mathbf{D}}_1 = \frac{\mathbf{d}_1 \dot{\gamma}_1 \sin(\theta_1 - \gamma_1)}{\sin \theta_1}$$





Figure 4. Velocity diagram of CC Engine

From the curve and their 6th order polynomial trendlines above, it was observed that at the extreme displacements, the velocities were not zero instead the zero mark appears at a very brief instance. This falls in line with the requirements for a smooth exhaust and intake, thus a smooth running of the engine.

2.5 Acceleration Calculations

With the information given above, a study of how the different core movable components accelerate at a supposed rotational speed of the rotor, that is, 3000 rev/minute, was carried out. From the loop closure equation, was deduced the following;

Equation 11:

 $\mathbf{R}_{1}\ddot{\mathbf{\theta}}_{1}(-\sin\theta_{1},\cos\theta_{1}) + \ddot{\mathbf{D}}_{1}(\cos 3\pi/2,\sin 3\pi/2) = \mathbf{R}_{1}\dot{\mathbf{\theta}}_{1}^{2}(\cos\theta_{1},\sin\theta_{1}) + \mathbf{d}_{1}\dot{\gamma}_{1}^{2}(\cos\gamma_{1},\sin\gamma_{1})$ Using the unit vectors quoted earlier, the scalar products with (11) deduced;

For $(\cos 3\pi/2, \sin 3\pi/2)$:

 $R_1\ddot{\theta}_1(-\sin\theta_1\cos 3\pi/2 + \cos\theta_1\sin 3\pi/2) + \ddot{D}_1$

 $= R_1 \dot{\theta}_1^2 (\cos \theta_1 \cos 3\pi/2 + \sin \theta_1 \sin 3\pi/2) + d_1 \dot{\gamma}_1^2 (\cos \gamma_1 \cos 3\pi/2 + \sin \gamma_1 \sin 3\pi/2)$

Equation 12:

$$\Rightarrow \ddot{D}_1 - R_1 \ddot{\theta}_1 \cos \theta_1 = -R_1 \dot{\theta}_1^2 \sin \theta_1 - d_1 \dot{\gamma}_1^2 \sin \gamma_1$$

And for $(-\sin 3\pi/2, \cos 3\pi/2)$

$$R_1 \ddot{\theta}_1 (\sin \theta_1 \sin 3\pi/2 + \cos \theta_1 \cos 3\pi/2)$$

= $R_1 \dot{\theta}_1^2 (-\cos \theta_1 \sin 3\pi/2 + \sin \theta_1 \cos 3\pi/2)$
+ $d_1 \dot{\gamma}_1^2 (-\cos \gamma_1 \sin 3\pi/2 + \sin \gamma_1 \cos 3\pi/2)$

Equation 13:

$$\Rightarrow - R_1 \ddot{\theta}_1 \sin \theta_1 = R_1 \dot{\theta}_1^2 \cos \theta_1 + d_1 \dot{\gamma}_1^2 \cos \gamma_1$$

From equations 12 and 13, we can result to; Equation 14:

$$\ddot{\theta}_1 = \frac{-R_1 \dot{\theta}_1^2 \cos \theta_1 - d_1 \dot{\gamma}_1^2 \cos \gamma_1}{R_1 \sin \theta_1}$$

And

$$\ddot{D}_{1} = R_{1}\ddot{\theta}_{1}\cos\theta_{1} - R_{1}\dot{\theta}_{1}^{2}\sin\theta_{1} - d_{1}\dot{\gamma}_{1}^{2}\sin\gamma_{1}$$

$$\Rightarrow \ddot{D}_{1} = R_{1}\cos\theta_{1}\left(\frac{-R_{1}\dot{\theta}_{1}^{2}\cos\theta_{1} - d_{1}\dot{\gamma}_{1}^{2}\cos\gamma_{1}}{R_{1}\sin\theta_{1}}\right) - R_{1}\dot{\theta}_{1}^{2}\sin\theta_{1} - d_{1}\dot{\gamma}_{1}^{2}\sin\gamma_{1}$$

$$\Rightarrow \ddot{D}_{1} = \frac{-R_{1}\dot{\theta}_{1}^{2}\cos^{2}\theta_{1} - d_{1}\dot{\gamma}_{1}^{2}\cos\gamma_{1}\cos\theta_{1} - R_{1}\dot{\theta}_{1}^{2}\sin^{2}\theta_{1} - d_{1}\dot{\gamma}_{1}^{2}\sin\gamma_{1}\sin\theta_{1}}{\sin\theta_{1}}$$

Equation 15:

$$\Rightarrow \ddot{\mathrm{D}}_{1} = \frac{-\mathrm{R}_{1}\dot{\theta}_{1}^{2} - \mathrm{d}_{1}\dot{\gamma}_{1}^{2}\cos(\gamma_{1} - \theta_{1})}{\sin\theta_{1}}$$

A graphical display of $\ddot{\theta}_1$ and \ddot{D}_1 can be perceived below;



Figure 5. Acceleration diagram of CC Engine

The above graphical presentation of \ddot{D}_1 its respective 6th order polynomial trendline proof that shortly before the intake, that is, when D_1 is at its maximum, the cam would have decelerated enough so that by the start of the expansion, there is very little or no resistance due to further displacement in the reactive direction and so an efficient use of the potential energy stored in the compressed air. $\ddot{\theta}_1$ shows a steady angular deceleration whose peak (when it touches the horizontal axis) is attained during the intake and exhaust. With our knowledge so far on the kinematics of the CC engine, we can see that the minimum pressure angle of our mechanism is zero.

Vividly these graphical representations brought in some satisfaction on the kinematics of the engine, to be better convinced with the functioning of CC engine, the dynamic formulations of the engine was done. (Psanis, 2007)

3. Internal Dynamics Formulations of the Canis Compressed Air (CC Engine) Engine

3.1 Skeleton of Dynamic Forces CC Engine

Consider figure 1, the links acting at the wing contact base of expansion chamber 1 (EC1) were isolated to give the following skeleton diagram (it was assumed the masses were insignificant).



Figure 6. Skeleton of CC engine with dynamic forces

3.2 Free Body Diagram of CC Engine

The free body diagram of the CC engine was established as can be seen below. For each link, the local unit vector tangent and normal were represented by ξ and η . R_{1D} represented an unknown resultant and M₁ a known torque; the expected torque of the engine.

3.3 Dynamic Formulation

Link d_1 ;

And

Thus

Equation 16:

$$F_{1d} = \frac{-M_1}{d_1 \sin(\theta_1 - \gamma_1)}$$

From the free body diagram, it was deduced that; Equation 17:

Link D_1 :

$$R_{1D} + T_{1D} + N_{1D} = 0$$

$$R_{1D}(\cos\theta_1, \sin\theta_1)^T + T_{1D}(-1, 0)^T + N_{1D}(0, -1)^T = 0$$

Thus

Equation 18:

$$T_{1D} = R_{1D}\cos\theta_1$$

And



Figure 7. Free body diagram of CC engine

With a rotor torque of 20 units, a graphical study of the forces was done as can be seen below.



Figure 8. Dynamic forces of CC Engine

The peak of the resultant force shortly after the intake and its trough at the exhaust brought about the conclusion that the system will react readily according to expectation, producing the maximum rotatory force per expansion chamber, during the intake. This will permit the mechanism to overcome forces of inertia. The mathematical analysis show that the resultant force gets to about ten times the average shortly after the intake and it drops down by similar figures during the exhaust.

4. Conclusion

This paper concludes with the study of the internal dynamic forces which are governed by the action – reaction principles. The findings at this level further concretize the susceptibility of a perfect intake, explosive expansion and smooth exhaust. With this much understanding of the system design, the smooth functioning of this engine is assured theoretically with little or no conflicting forces and movements. The next step will have to deal with the external forces which will yield the appropriate torque and speed of the engine thus a step closer to dimensioning the engine component parts.

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