



Existence of Positive Periodic Solutions of a Lotka-Volterra System with Multiple Time Delays

Changjin Xu

Faculty of Science, Hunan Institute of Engineering
Xiangtan 411004, China

Department of Mathematics, Central South University
Changsha 410083, China

Tel: 86-732-268-0745 E-mail: xcj403@126.com

Maoxin Liao

Department of Mathematics, Central South University
Changsha 410083, China

School of Mathematics and Physics, Nanhua University
Hengyang 421001, China

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Abstract

In this paper, a class of Lotka-Volterra system with multiple time delays is considered. By using the continuation theorem of coincidence degree theory, we derive a set of easily verifiable sufficient conditions that guarantees the existence of at least a positive periodic solution.

Keywords: Lotka-Volterra system, Periodic solution, Multiple time delay, Continuation theorem, Topological degree

1. Introduction

Lotka-Volterra system is an important population system and has been studied by many authors, see (Chen Shihua, et.al, 2004, Fan Meng, et.al, 1999, Li Yongkun, et.al, 2001) and the reference therein. but most of the previous results focused on the stability, attractiveness, persistence and periodicity of solution to the ordinary differential systems or time delay systems with constant delays. Rare work has been done for the systems with varying delays and varying coefficients. In 1991, Weng di Wang et. al (Wang Wending, et. al, 1991) had considered a two-dimensional predator-prey system with a finite constant number discrete delays

$$\begin{cases} \dot{x}(t) = x(t) \left[r_1 - \sum_{j=1}^m a_{1j} x(t - \tau_{1j}) - \sum_{j=1}^m b_{1j} y(t - \rho_{1j}) \right], \\ \dot{y}(t) = y(t) \left[r_2 + \sum_{j=1}^m a_{2j} x(t - \tau_{2j}) - \sum_{j=1}^m b_{2j} y(t - \rho_{2j}) \right], \end{cases} \quad (1)$$

with initial conditions

$$x(s) = \varphi(s) \geq 0, s \in [-\tau, 0]; \varphi(0) > 0,$$

$$y(s) = \psi(s) \geq 0, s \in [-\tau, 0]; \psi(0) > 0,$$

where r_1, r_2 are real constants with $r_1 > 0$; $a_{ij}, b_{ij}, \tau_{ij}, \rho_{ij}$ ($i = 1, 2; j = 1, 2, \dots, m$) are non-negative constants. Not all of a_{1j} and not all b_{1j} ($j = 1, 2, \dots, m$) are zero; Both $\varphi(s)$ and $\psi(s)$ are continuous on the interval $[-\tau, 0]$ in which $\tau = \max\{\tau_{ij}, \rho_{ij} : i = 1, 2; j = 1, 2, \dots, m\}$. And obtained the conclusion that the time delays are harmless for uniform persistence of the solutions to the system.

We note that any biological or environmental parameters are naturally subject to fluctuation in time. It is necessary and important to consider models with periodic ecological parameters or perturbations which might be naturally exposed (for example, those due to seasonal effects of weather, food supply, mating habits, hunting or harvesting seasons, etc.). Thus, the assumption of periodicity of the parameters is a way of incorporating the periodicity of the environment.

In this paper, we are concerned with the effects of periodicity of ecological and environmental parameters and time delays. Then system (1) can be modified as the form:

$$\begin{cases} \dot{x}(t) = x(t) \left[r_1(t) - \sum_{j=1}^m a_{1j}(t)x(t - \tau_{1j}(t)) - \sum_{j=1}^m b_{1j}(t)y(t - \rho_{1j}(t)) \right], \\ \dot{y}(t) = y(t) \left[r_2(t) + \sum_{j=1}^m a_{2j}(t)x(t - \tau_{2j}(t)) - \sum_{j=1}^m b_{2j}(t)y(t - \rho_{2j}(t)) \right], \end{cases} \quad (2)$$

with initial conditions

$$\begin{cases} x(s) = \varphi(s) \geq 0, s \in [-\tau, 0]; \varphi(0) > 0, \\ y(s) = \psi(s) \geq 0, s \in [-\tau, 0]; \psi(0) > 0, \end{cases} \quad (3)$$

where $r_1(t), r_2(t)$ are real functions with $r_i(t) > 0, (i = 1, 2)$; $a_{ij}(t), b_{ij}(t), \tau_{ij}(t), \rho_{ij}(t) (i = 1, 2; j = 1, 2, \dots, m)$ are non-negative functions. Not all of $a_{1j}(t)$ and not all $b_{1j}(t) (j = 1, 2, \dots, m)$ are zero; Both $\varphi(s)$ and $\psi(s)$ are continuous on the interval $[-\tau, 0]$ in which $\tau = \max_{t \in R} \max\{\tau_{ij}(t), \rho_{ij}(t) : i = 1, 2; j = 1, 2, \dots, m\}$.

Throughout the paper, we always assume that

(H₁) $r_i(t), a_{ij}(t), b_{ij}(t), \tau_{ij}(t), \rho_{ij}(t) (i = 1, 2; j = 1, 2, \dots, m)$ are ω periodic, i.e.,

$$r_i(t + \omega) = r_i(t), a_{ij}(t + \omega) = a_{ij}(t), b_{ij}(t + \omega) = b_{ij}(t),$$

$$\tau_{ij}(t + \omega) = \tau_{ij}(t), \rho_{ij}(t + \omega) = \rho_{ij}(t)$$

for any $t \in R$.

(H₂) $r_i(t), a_{ij}(t), b_{ij}(t), \tau_{ij}(t), \rho_{ij}(t) (i = 1, 2; j = 1, 2, \dots, m)$ are all positive, i.e.,

$$r_i(t), a_{ij}(t), b_{ij}(t), \tau_{ij}(t), \rho_{ij}(t) (i = 1, 2; j = 1, 2, \dots, m) > 0.$$

The principle object of this article is to find a set of sufficient conditions that guarantees the existence of at least a positive periodic solution for system (2) (3).

2. Basic lemma

In order to explore the existence of positive periodic solutions of (2) (3) and for the reader's convenience, we shall first summarize below a few concepts and results without proof, borrowing from (Yang Zhihui, et. al, 2007).

Let X, Y be normed vector spaces, $L : \text{Dom} L \subset X \rightarrow Y$ is a linear mapping, $N : X \rightarrow Y$ is a continuous mapping. The mapping L will be called a Fredholm mapping of index zero if $\dim \text{Ker} L = \text{codim Im} L < +\infty$ and $\text{Im} L$ is closed in Y . If L is a Fredholm mapping of index zero and there exist continuous projectors $P : X \rightarrow X$ and $Q : Y \rightarrow Y$ such that $\text{Im} P = \text{Ker} L, \text{Im} L = \text{Ker} Q = \text{Im}(I - Q)$, it follows that $L|_{\text{Dom} L \cap \text{Ker} P} : (I - P)X \rightarrow \text{Im} L$ is invertible. We denote the inverse of that map by K_P . If Ω is an open bounded subset of X , the mapping N will be called L -compact on $\bar{\Omega}$ if $QN(\bar{\Omega})$ is bounded and $K_P(I - Q)N : \bar{\Omega} \rightarrow X$ is compact. Since $\text{Im} Q$ is isomorphic to $\text{Ker} L$, there exist isomorphisms $J : \text{Im} Q \rightarrow \text{Ker} L$.

Lemma 2.1. (Robert E. Gaines et. al, 1991)(Continuation Theorem) Let L be a fredholm mapping of index zero and let N be L -compact on $\bar{\Omega}$. Suppose

- (a) for each $\lambda \in (0, 1)$, every solution x of $Lx = \lambda Nx$ is such that $x \notin \partial\Omega$;
 - (b) $QNx \neq 0$ for each $x \in \text{Ker} L \cap \partial\Omega$, and $\deg\{JQN, \Omega \cap \partial\text{Ker} L, 0\} \neq 0$;
- Then the equation $Lx = Nx$ has at least one solution lying in $\text{Dom} L \cap \bar{\Omega}$.

Lemma 2.2. $R_+^2 = \{(x(t), y(t))^T \in R^2 \mid x(t) > 0, y(t) > 0\}$ is positive invariant with respect to system (2) (3).

Proof. In fact,

$$\begin{aligned} x(t) &= \varphi(0) \exp \int_0^t \left[r_1(s) - \sum_{j=1}^m a_{1j}(s)x(s - \tau_{1j}(s)) - \sum_{j=1}^m b_{1j}(s)y(s - \rho_{1j}(s)) \right] ds, \\ y(t) &= \psi(0) \exp \int_0^t \left[r_2(s) + \sum_{j=1}^m a_{2j}(s)x(s - \tau_{2j}(s)) - \sum_{j=1}^m b_{2j}(s)y(s - \rho_{2j}(s)) \right] ds. \end{aligned}$$

In view of $\varphi(0) > 0, \psi(0) > 0, (i = 1, 2)$, obviously, the conclusion follows.

3. Existence of positive periodic solutions

For convenience and simplicity in the following discussion, we always use the notations below throughout the paper:

$$\bar{g} = \frac{1}{\omega} \int_0^\omega g(t) dt, \quad g^L = \min_{t \in [0, \omega]} g(t), \quad g^M = \max_{t \in [0, \omega]} g(t),$$

where $g(t)$ is a ω continuous periodic function. Let $\sigma_{ij}(t) = t - \tau_{ij}(t)$, $\theta_{ij}(t) = t - \rho_{ij}(t)$, $t \in R$, $i = 1, 2$; $j = 1, 2, \dots, m$. Assume that

$$(H_3) \quad \tau'_{ij}(t) < 1, \quad \rho'_{ij}(t) < 1, \quad (i = 1, 2; j = 1, 2, \dots, m)$$

Then $\sigma_{ij}(t)$ and $\theta_{ij}(t)$ have inverse functions denoted by $\mu_{ij}(t)$, $\varepsilon_{ij}(t)$, ($i = 1, 2$; $j = 1, 2, \dots, m$), respectively. Obviously, $\mu_{ij}(t + \omega) = \mu_{ij}(t) + \omega$, $\varepsilon_{ij}(t + \omega) = \varepsilon_{ij}(t) + \omega$.

In the following, we will ready to state and prove our result.

Theorem 3.1. Suppose that (H_1) , (H_2) , (H_3) , (H_4) $D_{22}^L \bar{r}_1^L > D_{12}^M \bar{r}_2^M$ and

$$(H_5) \quad \bar{r}_1 \sum_{j=1}^m \bar{b}_{2j} > \bar{r}_2 \sum_{j=1}^m \bar{b}_{1j} \text{ hold, where}$$

$$D_{22} = \sum_{j=1}^m \frac{b_{2j}(t) \varepsilon_{2j}(t)}{1 - \rho'_{2j}(\varepsilon_{2j}(t))}, \quad D_{12} = \sum_{j=1}^m \frac{b_{1j}(t) \varepsilon_{1j}(t)}{1 - \rho'_{1j}(\varepsilon_{1j}(t))},$$

then the system (2) (3) has at least a ω periodic solution.

Proof. Since solutions of (2) (3) remained positive for all $t \geq 0$, we let

$$u_1(t) = \ln[x(t)], \quad u_2(t) = \ln[y(t)]. \quad (4)$$

Substituting (4) into (2), we obtain

$$\begin{cases} \dot{u}_1(t) = r_1(t) - \sum_{j=1}^m a_{1j}(t) \exp\{u_1(t - \tau_{1j}(t))\} - \sum_{j=1}^m b_{1j}(t) \exp\{u_2(t - \rho_{1j}(t))\}, \\ \dot{u}_2(t) = r_2(t) + \sum_{j=1}^m a_{2j}(t) \exp\{u_1(t - \tau_{2j}(t))\} - \sum_{j=1}^m b_{2j}(t) \exp\{u_2(t - \rho_{2j}(t))\}. \end{cases} \quad (5)$$

It is easy to see that if system (5) has one ω periodic solution $(u_1^*(t), u_2^*(t))^T$, then $(x^*(t), y^*(t))^T = (\exp[u_1^*(t)], \exp[u_2^*(t)])^T$ is a positive solution of system (2). Therefore, to complete the proof, it suffices to show that system (5) has at least one ω periodic solution.

Let $X = Z = \{u(t)\} = \{(u_1(t), u_2(t))^T \mid u(t) \in C(R, R^2), u(t + \omega) = u(t)\}$, and define $\|u\| = \|(u_1(t), u_2(t))^T\| = \max_{t \in [0, \omega]} |u_1(t)| + \max_{t \in [0, \omega]} |u_2(t)|$. Then X and Z are Banach spaces when they are endowed with the norm $\|\cdot\|$.

Let $L : \text{Dom} L \subset X \rightarrow Z$ and $N : X \rightarrow Z$ be the following:

$$\begin{aligned} Lu &= \dot{x}'(t), \\ Nu &= \begin{pmatrix} r_1(t) - \sum_{j=1}^m a_{1j}(t) \exp\{u_1(t - \tau_{1j}(t))\} - \sum_{j=1}^m b_{1j}(t) \exp\{u_2(t - \rho_{1j}(t))\} \\ r_2(t) + \sum_{j=1}^m a_{2j}(t) \exp\{u_1(t - \tau_{2j}(t))\} - \sum_{j=1}^m b_{2j}(t) \exp\{u_2(t - \rho_{2j}(t))\} \end{pmatrix}. \end{aligned} \quad (6)$$

Define continuous projective operators P and Q :

$$Pu = \frac{1}{\omega} \int_0^\omega u(t) dt, \quad Qu = \frac{1}{\omega} \int_0^\omega u(t) dt, \quad u \in X, u \in Z.$$

We can see that $\text{Ker} L = \{u \in X \mid u = h \in R^2\}$, $\text{Im} L = \{u \in Z \mid \int_0^\omega u(t) dt = 0\}$ is closed in X and $\dim(\text{Ker} L) = 2 = \text{codim}(\text{Im} L)$, then it follows that L is a Fredholm mapping of index zero. Moreover, it is easy to check that

$$QNx = \begin{pmatrix} \frac{1}{\omega} \int_0^\omega \left[r_1(t) - \sum_{j=1}^m a_{1j}(t) \exp\{u_1(t - \tau_{1j}(t))\} - \sum_{j=1}^m b_{1j}(t) \exp\{u_2(t - \rho_{1j}(t))\} \right] dt \\ \frac{1}{\omega} \int_0^\omega \left[r_2(t) + \sum_{j=1}^m a_{2j}(t) \exp\{u_1(t - \tau_{2j}(t))\} - \sum_{j=1}^m b_{2j}(t) \exp\{u_2(t - \rho_{2j}(t))\} \right] dt \end{pmatrix}.$$

By easily computation, we have

$$K_P(z) = \int_0^\omega z(u)du - \frac{1}{\omega} \int_0^\omega \left[\int_0^t z(u)du \right] dt,$$

$$K_P(I - Q)Nu =$$

$$\begin{pmatrix} \int_0^t F_1(s)ds \\ \int_0^t F_2(s)ds \end{pmatrix} - \begin{pmatrix} \frac{1}{\omega} \int_0^\omega \int_0^t F_1(s)dsdt \\ \frac{1}{\omega} \int_0^\omega \int_0^t F_2(s)dsdt \end{pmatrix} - \begin{pmatrix} (\frac{t}{\omega} - \frac{1}{2}) \int_0^\omega F_1(s)ds \\ (\frac{t}{\omega} - \frac{1}{2}) \int_0^\omega F_2(s)ds \end{pmatrix}, \quad (7)$$

where

$$\begin{aligned} F_1(s) &= r_1(s) - \sum_{j=1}^m a_{1j}(s) \exp\{u_1(s - \tau_{1j}(s))\} - \sum_{j=1}^m b_{1j}(s) \exp\{u_2(s - \rho_{1j}(s))\}, \\ F_2(s) &= r_2(s) + \sum_{j=1}^m a_{2j}(s) \exp\{u_1(s - \tau_{2j}(s))\} - \sum_{j=1}^m b_{2j}(s) \exp\{u_2(s - \rho_{2j}(s))\}. \end{aligned}$$

Obviously, QN and $K_P(I - Q)N$ are continuous. Since X is a finite-dimensional Banach space, using the Ascoli-Arzelà theorem, it is not difficult to show that $K_P(I - Q)N(\bar{\Omega})$ is compact for any open bounded set $\Omega \subset X$. Moreover, $QN(\bar{\Omega})$ is bounded. Thus, N is L -compact on $\bar{\Omega}$ with any open bounded set $\Omega \subset X$.

Now we are at the point to search for an appropriate open, bounded subset Ω for the application of the continuation theorem. Corresponding to the operator equation $Ly = \lambda Ny$, $\lambda \in (0, 1)$, we have

$$\begin{cases} u_1(t) = \lambda \left[r_1(t) - \sum_{j=1}^m a_{1j}(t) \exp\{u_1(t - \tau_{1j}(t))\} - \sum_{j=1}^m b_{1j}(t) \exp\{u_2(t - \rho_{1j}(t))\} \right], \\ u_2(t) = \lambda \left[r_2(t) + \sum_{j=1}^m a_{2j}(t) \exp\{u_1(t - \tau_{2j}(t))\} - \sum_{j=1}^m b_{2j}(t) \exp\{u_2(t - \rho_{2j}(t))\} \right]. \end{cases} \quad (8)$$

Suppose that $u(t) = (u_1(t), u_2(t))^T \in X$ is an arbitrary solution of system (8) for a certain $\lambda \in (0, 1)$, integrating both sides of (8) over the interval $[0, \omega]$ with respect to t , we obtain

$$\begin{cases} \int_0^\omega \left[\sum_{j=1}^m a_{1j}(t) \exp\{u_1(t - \tau_{1j}(t))\} + \sum_{j=1}^m b_{1j}(t) \exp\{u_2(t - \rho_{1j}(t))\} \right] dt = \bar{r}_1 \omega, \\ \int_0^\omega \left[\sum_{j=1}^m a_{2j}(t) \exp\{u_1(t - \tau_{2j}(t))\} - \sum_{j=1}^m b_{2j}(t) \exp\{u_2(t - \rho_{2j}(t))\} \right] dt = -\bar{r}_2 \omega. \end{cases} \quad (9)$$

In view of the following:

$$\begin{aligned} & \int_0^\omega \sum_{j=1}^m a_{ij}(t) \exp\{u_1(t - \tau_{ij}(t))\} dt \\ &= \int_{-\tau_{ij}(0)}^{\omega - \tau_{ij}(\omega)} \sum_{j=1}^m \frac{a_{ij}(\mu_{ij}(s))}{1 - \tau'_{ij}(\mu_{ij}(s))} \exp\{u_1(s)\} ds, \\ &= \int_{-\tau_{ij}(0)}^{\omega - \tau_{ij}(\omega)} \sum_{j=1}^m \frac{a_{ij}(\mu_{ij}(s))}{1 - \tau'_{ij}(\mu_{ij}(s))} \exp\{u_1(s)\} ds \\ &= \int_0^\omega \sum_{j=1}^m \frac{a_{ij}(\mu_{ij}(s))}{1 - \tau'_{ij}(\mu_{ij}(s))} \exp\{u_1(s)\} ds, (i = 1, 2; j = 1, 2, \dots, m), \end{aligned} \quad (10)$$

$$\begin{aligned} & \int_0^\omega \sum_{j=1}^m b_{ij}(t) \exp\{u_2(t - \rho_{ij}(t))\} dt \\ &= \int_{-\rho_{ij}(0)}^{\omega - \rho_{ij}(\omega)} \sum_{j=1}^m \frac{b_{ij}(\varepsilon_{ij}(s))}{1 - \rho'_{ij}(\varepsilon_{ij}(s))} \exp\{u_2(s)\} ds \\ &= \int_{-\rho_{ij}(0)}^{\omega - \rho_{ij}(\omega)} \sum_{j=1}^m \frac{b_{ij}(\varepsilon_{ij}(s))}{1 - \rho'_{ij}(\varepsilon_{ij}(s))} \exp\{u_2(s)\} ds \\ &= \int_0^\omega \sum_{j=1}^m \frac{b_{ij}(\varepsilon_{ij}(s))}{1 - \rho'_{ij}(\varepsilon_{ij}(s))} \exp\{u_2(s)\} ds, (i = 1, 2; j = 1, 2, \dots, m). \end{aligned} \quad (11)$$

From (9) (10) (11), we can obtain

$$\int_0^\omega \sum_{j=1}^m \frac{a_{1j}(\mu_{1j}(s))}{1 - \tau'_{1j}(\mu_{1j}(s))} \exp\{u_1(s)\} ds + \int_0^\omega \sum_{j=1}^m \frac{b_{1j}(\varepsilon_{1j}(s))}{1 - \rho'_{1j}(\varepsilon_{1j}(s))} \exp\{u_2(s)\} ds = \bar{r}_1 \omega, \quad (12)$$

$$\int_0^\omega \sum_{j=1}^m \frac{a_{2j}(\mu_{2j}(s))}{1 - \tau'_{2j}(\mu_{2j}(s))} \exp\{u_1(s)\} ds - \int_0^\omega \sum_{j=1}^m \frac{b_{2j}(\varepsilon_{2j}(s))}{1 - \rho'_{2j}(\varepsilon_{2j}(s))} \exp\{u_2(s)\} ds = -\bar{r}_2 \omega. \quad (13)$$

By the mean value theorem for improper integral, there exist $\xi_{ikj} \in [0, \omega]$ ($i = 1, 2; k = 1, 2; j = 1, 2, \dots, m$) such that

$$A_{11} \int_0^\omega \exp\{u_1(s)\} ds + A_{12} \int_0^\omega \exp\{u_2(s)\} ds = \bar{r}_1 \omega, \quad (14)$$

$$A_{21} \int_0^\omega \exp\{u_1(s)\} ds - A_{22} \int_0^\omega \exp\{u_2(s)\} ds = -\bar{r}_2 \omega, \quad (15)$$

where

$$A_{11} = \sum_{j=1}^m \frac{a_{1j}(\xi_{11j})\mu_{1j}(\xi_{11j})}{1 - \tau'_{1j}(\mu_{1j}(\xi_{11j}))},$$

$$A_{12} = \sum_{j=1}^m \frac{b_{1j}(\xi_{12j})\varepsilon_{1j}(\xi_{12j})}{1 - \rho'_{1j}(\varepsilon_{1j}(\xi_{12j}))},$$

$$A_{21} = \sum_{j=1}^m \frac{a_{2j}(\xi_{21j})\mu_{2j}(\xi_{21j})}{1 - \tau'_{2j}(\mu_{2j}(\xi_{21j}))},$$

$$A_{22} = \sum_{j=1}^m \frac{b_{2j}(\xi_{22j})\varepsilon_{2j}(\xi_{22j})}{1 - \rho'_{2j}(\varepsilon_{2j}(\xi_{22j}))}.$$

Then,

$$\int_0^\omega \exp\{u_1(s)\} ds = \frac{(A_{22}\bar{r}_1 - A_{12}\bar{r}_2)\omega}{A_{11}A_{22} + A_{21}A_{12}}, \quad (16)$$

$$\int_0^\omega \exp\{u_2(s)\} ds = \frac{(A_{21}\bar{r}_1 + A_{11}\bar{r}_2)\omega}{A_{11}A_{22} + A_{21}A_{12}}. \quad (17)$$

So we have

$$\frac{(D_{22}^L \bar{r}_1^L - D_{12}^M \bar{r}_2^M)\omega}{D_{11}^M D_{22}^M + D_{21}^M D_{12}^M} \leq \int_0^\omega \exp\{u_1(s)\} ds \leq \frac{(D_{22}^M \bar{r}_1^M - D_{12}^L \bar{r}_2^L)\omega}{D_{11}^L D_{22}^L + D_{21}^L D_{12}^L}, \quad (18)$$

$$\frac{(D_{21}^L \bar{r}_1^L + D_{11}^M \bar{r}_2^M)\omega}{D_{11}^M D_{22}^M + D_{21}^M D_{12}^M} \leq \int_0^\omega \exp\{u_2(s)\} ds \leq \frac{(D_{21}^M \bar{r}_1^M + D_{11}^L \bar{r}_2^L)\omega}{D_{11}^L D_{22}^L + D_{21}^L D_{12}^L}, \quad (19)$$

where

$$D_{11} = \sum_{j=1}^m \frac{a_{1j}(t)\mu_{1j}(t)}{1 - \tau'_{1j}(\mu_{1j}(t))},$$

$$D_{12} = \sum_{j=1}^m \frac{b_{1j}(t)\varepsilon_{1j}(t)}{1 - \rho'_{1j}(\varepsilon_{1j}(t))},$$

$$D_{21} = \sum_{j=1}^m \frac{a_{2j}(t)\mu_{2j}(t)}{1 - \tau'_{2j}(\mu_{2j}(t))},$$

$$D_{22} = \sum_{j=1}^m \frac{b_{2j}(t)\varepsilon_{2j}(t)}{1 - \rho'_{2j}(\varepsilon_{2j}(t))}.$$

By the condition of theorem 3.1, there exist $t_i \in [0, \omega], i = 1, 2$ such that

$$u_1(t_1) = \ln \left[\frac{(A_{22}\bar{r}_1 - A_{12}\bar{r}_2)\omega}{A_{11}A_{22} + A_{21}A_{12}} \right], \quad (20)$$

$$u_2(t_2) = \ln \left[\frac{(A_{21}\bar{r}_1 + A_{11}\bar{r}_2)\omega}{A_{11}A_{22} + A_{21}A_{12}} \right]. \quad (21)$$

By the condition (H_4) of theorem 3.1, there exist $B_1, B_2 > 0$ such that

$$|u_1(t_1)| \leq B_1, |u_2(t_2)| \leq B_2. \quad (22)$$

In view of the following:

$$\begin{aligned} \int_0^\omega |\dot{u}_1(t)| dt &= \lambda \int_0^\omega \left\| r_1(t) - \sum_{j=1}^m a_{1j}(t) \exp\{u_1(t - \tau_{1j}(t))\} \right. \\ &\quad \left. - \sum_{j=1}^m b_{1j}(t) \exp\{u_2(t - \rho_{1j}(t))\} \right\| dt \\ &\leq 2 \int_0^\omega |r_1(t)| dt = 2 \int_0^\omega r_1(t) dt = 2\bar{r}_1\omega := B_3, \end{aligned} \quad (23)$$

$$\begin{aligned} \int_0^\omega |\dot{u}_2(t)| dt &= \lambda \int_0^\omega \left\| r_2(t) + \sum_{j=1}^m a_{2j}(t) \exp\{u_1(t - \tau_{2j}(t))\} \right. \\ &\quad \left. - \sum_{j=1}^m b_{2j}(t) \exp\{u_2(t - \rho_{2j}(t))\} \right\| dt \\ &\leq \bar{r}_2\omega + \int_0^\omega \left| \sum_{j=1}^m a_{2j}(t) \exp\{u_1(t - \tau_{2j}(t))\} \right| dt \\ &\quad + \int_0^\omega \left| \sum_{j=1}^m b_{2j}(t) \exp\{u_2(t - \rho_{2j}(t))\} \right| dt \\ &\leq \bar{r}_2\omega + \sum_{j=1}^m \frac{a_{2j}(\mu_{2j}(s))\omega}{1 - \mu'_{2j}(\mu_{2j}(s))} \int_0^\omega \exp\{u_1(s)\} ds \\ &\quad + \sum_{j=1}^m \frac{b_{2j}(\varepsilon_{2j}(s))\omega}{1 - \rho'_{2j}(\varepsilon_{2j}(s))} \int_0^\omega \exp\{u_2(s)\} ds. \\ &\leq \bar{r}_2\omega + \sum_{j=1}^m \frac{a_{2j}(\mu_{2j}(s))\omega}{1 - \mu'_{2j}(\mu_{2j}(s))} \frac{(A_{22}^M \bar{r}_1^M - A_{12}^L \bar{r}_2^L)\omega}{A_{11}^L A_{22}^L + A_{21}^L A_{12}^L} \\ &\quad + \sum_{j=1}^m \frac{b_{2j}(\varepsilon_{2j}(s))\omega}{1 - \rho'_{2j}(\varepsilon_{2j}(s))} \frac{(A_{21}^M \bar{r}_1^M + A_{11}^L \bar{r}_2^L)\omega}{A_{11}^L A_{22}^L + A_{21}^L A_{12}^L} := B_4, \end{aligned} \quad (24)$$

then it follows from (22) (23) (24) that

$$|u_1(t)| \leq |u_1(t_1)| + \int_0^\omega |\dot{u}_1(t)| dt \leq B_1 + B_3 := B_5, \quad (25)$$

$$|u_2(t)| \leq |u_2(t_2)| + \int_0^\omega |\dot{u}_2(t)| dt \leq B_2 + B_4 := B_6. \quad (26)$$

Obviously, B_1, B_2, B_3, B_4 are independent of $\lambda \in (0, 1)$. By the condition (H_5) of Theorem 3.1, it is easy to show that the algebraic equations

$$\begin{cases} \sum_{i=1}^m \bar{a}_{1j} u_1 + \sum_{i=1}^m \bar{b}_{1j} u_2 = \bar{r}_1, \\ \sum_{i=1}^m \bar{a}_{2j} u_1 - \sum_{i=1}^m \bar{b}_{2j} u_2 = -\bar{r}_2, \end{cases} \quad (27)$$

has a unique positive solution $(\tilde{u}_1, \tilde{u}_2)^T \in R^2$. Take $M = \max\{B_5, B_6\} + B_0$, where B_0 is taken sufficiently large such that the unique positive solution $(\tilde{u}_1, \tilde{u}_2)^T \in R^2$ satisfies $\max_{t \in [0, \omega]} |\tilde{u}_1| + \max_{t \in [0, \omega]} |\tilde{u}_2| < B_0$,

Let $\Omega := \{u = \{u(t)\} \in X : \|u\| < M\}$, then it is easy to see that Ω is an open, bounded set in X and verifies requirement (a) of Lemma 2.1. When $(u_1(t), u_2(t))^T \in \partial\Omega \cap \text{Ker} L = \partial\Omega \cap R^2$, $u = \{(u_1, u_2)^T\}$ is a constant vector in R^2 with $\|u\| =$

$\|(u_1(t), u_2(t))^T\| = \max_{t \in [0, \omega]} |u_1(t)| + \max_{t \in [0, \omega]} |u_2(t)| = M$. Then we have

$$QNy = \begin{pmatrix} \bar{r}_1 - \sum_{i=1}^m \bar{a}_{1j} u_1 - \sum_{i=1}^m \bar{b}_{1j} u_2 \\ \bar{r}_2 + \sum_{i=1}^m \bar{a}_{2j} u_1 - \sum_{i=1}^m \bar{b}_{2j} u_2 \end{pmatrix} \neq 0. \quad (28)$$

Letting J be the identity mapping and by direct calculation, we get

$$\begin{aligned} & \deg \{JQN(u_1, u_2)^T; \partial\Omega \cap \ker L; 0\} \\ &= \deg \{QN(u_1, u_2)^T; \partial\Omega \cap \ker L; 0\} \\ &= \text{sign} \left\{ \det \begin{pmatrix} -\sum_{i=1}^m \bar{a}_{1j} & -\sum_{i=1}^m \bar{b}_{1j} \\ \sum_{i=1}^m \bar{a}_{2j} & -\sum_{i=1}^m \bar{b}_{2j} \end{pmatrix} \right\} \\ &= \text{sign} \left\{ \sum_{i=1}^m \bar{a}_{1j} \sum_{i=1}^m \bar{b}_{2j} + \sum_{i=1}^m \bar{a}_{2j} \sum_{i=1}^m \bar{b}_{1j} \right\} = 1 \neq 0. \end{aligned}$$

This proves that condition (b) in Lemma 2.1 is satisfied. By now, we have proved that Ω verifies all requirements of Lemma 2.1, then it follows that $Lu = Nu$ has at least one solution $(u_1(t), u_2(t))^T$ in $\text{Dom} L \cap \bar{\Omega}$, that is to say, (5) has at least one ω periodic solution in $\text{Dom} L \cap \bar{\Omega}$. Then we know that $((x(t), y(t))^T = (\exp\{u_1(t)\}, \exp\{u_2(t)\})^T$ is an ω periodic solution of system (2) (3) with strictly positive components. This completes the proof.

Remark 3.1. Theorem 3.1 remains valid if some or all terms are replaced by corresponding terms with discrete time delays, distribute delays (finite or infinite), or deviating arguments respectively. At this point, we would like to point out that, when one applies the continuation theorem from the coincidence degree theory to explore the existence of periodic solutions to the system of differential equations or difference equations, time delays of any type or the deviating arguments have no effect on the existence of positive solutions.

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