# On the Irreducibility of Artin's Group of Graphs 

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Received: March 12, 2015 Accepted: March 30, 2015 Online Published: May 1, 2015
doi:10.5539/jmr.v7n2p117 URL: http://dx.doi.org/10.5539/jmr.v7n2p117


#### Abstract

We consider the graph $E_{3,1}$ with three generators $\sigma_{1}, \sigma_{2}, \delta$, where $\sigma_{1}$ has an edge with each of $\sigma_{2}$ and $\delta$. We then define the Artin group of the graph $E_{3,1}$ and consider its reduced Perron representation of degree three. After we specialize the indeterminates used in defining the representation to non-zero complex numbers, we obtain a necessary and sufficient condition that guarantees the irreducibility of the representation.


Keywords: Artin representation, braid group, Burau representation, graph, irreducibility

## 1. Introduction

To any undirected simple graph $T$, we introduce the Artin group, $A$, which is defined as an abstract group with vertices of $\Gamma$ as its generators and two relations: $x y=y x$ for vertices x and y that have no edge in common and $x y x=y x y$ if the vertices x and y have a common edge.
Let $A_{n}$ be the graph having $n$ vertices $\sigma_{i}$ 's $(1 \leq i \leq n)$ in which $\sigma_{i}$ and $\sigma_{i+1}$ share a comon edge, where $i=1,2, \ldots, n-1$. We notice that the Artin group of $A_{n}$ is the braid group on $n+1$ strands. That is, $A\left(A_{n}\right)=B_{n+1}$ (J.S.Birman, 1975).

Having defined $A_{n}$, we consider $E_{n+1, p}$, which is the graph obtained from $A_{n}$ by adding a vertex $\delta$ and an edge connecting $\sigma_{p}$ and $\delta$. Here $1 \leq p \leq n$. It is easy to see that the graph $A_{n}$ embeds in the graph $E_{n+1, p}$. That is, $A\left(A_{n}\right) \subset A\left(E_{n+1, p}\right)$. This induces an injection on $B_{n+1}$ to $A\left(E_{n+1, p}\right)$. In other words, a representation of $A\left(E_{n+1, p}\right)$ yields a representation of $B_{n+1}$.
Knowing the reduced Burau representation of $B_{n+1}$ of degree $n$, Perron extends such a representation to a representation of $B_{n+1}$ of degree $2 n$. The representation obtained is referred to as Burau bis representation. Next, Perron constructs for each $\lambda=\left(\lambda_{1}, \ldots, \lambda_{n}\right)$ a representation $\psi_{\lambda}: A\left(E_{n+1, p}\right) \rightarrow G L_{2 n}\left(Q\left(t, d_{1}, \ldots, d_{n}\right)\right)$, where $t, d_{1}, \ldots, d_{n} \lambda_{1}, \ldots, \lambda_{n}$ are indeterminates. We specialize $t, d_{1}, \ldots, d_{n}$ to non zero complex numbers, and we study this representation explicitly in the case $n=2$ and $p=1$. We then reduce the complex specialization of the representation $\psi_{\lambda}$ to a representation of degree 3 , namely $A\left(E_{3,1}\right) \rightarrow G L_{3}(\mathbb{C})$. A necessary and sufficient condition which guarantees its irreducibility is obtained in that case.

## 2. Burau bis Representation

Perron's strategy is to begin with the Burau representation of the braid group and extend it to a representation of $A\left(E_{n+1, p}\right)$. He begins with the reduced Burau representation: $B_{n+1} \rightarrow G L_{n}\left(\mathbb{Z}\left[t, t^{-1}\right]\right)$ defined as follows:

$$
\sigma_{i} \rightarrow J_{i}=\left(\begin{array}{c|ccc|c}
I_{i-2} & & 0 & & 0 \\
\hline & 1 & 0 & 0 & \\
0 & t & -t & 1 & 0 \\
& 0 & 0 & 1 & \\
\hline 0 & & 0 & & I_{n-i-1}
\end{array}\right)
$$

where $I_{k}$ stands for the $k \times k$ identity matrix. Here, $i=2, \ldots, n-1$.

$$
\begin{aligned}
& \sigma_{1} \rightarrow J_{1}=\left(\begin{array}{cc|c}
-t & 1 & 0 \\
0 & 1 & 0 \\
\hline 0 & I_{n-2}
\end{array}\right) \\
& \sigma_{n} \rightarrow J_{n}=\left(\begin{array}{c|cc}
I_{n-2} & 0 \\
\hline 0 & \begin{array}{c}
1 \\
0 \\
t
\end{array} \\
\hline
\end{array}\right)
\end{aligned}
$$

Knowing that this representation is of degree $n$, Perron extends it to a representation of $B_{n+1}$ of degree $2 n$. Let $R_{i}$ denote an $n \times n$ block of zeros with a $t$ placed in the $(i, i) t h$ position, and let $I_{n}$ denote the $n \times n$ identity matrix. The obtained representation is referred to as the Burau bis representation. It is defined as follows:

$$
\begin{gathered}
\psi: B_{n+1} \rightarrow G l_{2 n}\left(\mathbb{Z}\left[t, t^{-1}\right]\right) \\
\psi\left(\sigma_{i}\right)=\left(\begin{array}{ll}
I_{n} & 0 \\
R_{i} & J_{i}
\end{array}\right), \quad 1 \leq i \leq n
\end{gathered}
$$

For more details, see (T.E.Brendle, 2002, B.Perron, 1999).

## 3. Perron Representation

The Burau bis representation extends to $A\left(E_{n+1, p}\right)$ for all possible values of $n$ and $p$ in the following way.
Let $b=\left(\begin{array}{c}b_{1} \\ \vdots \\ b_{n}\end{array}\right), d=\left(\begin{array}{c}d_{1} \\ \vdots \\ d_{n}\end{array}\right)$, and $\lambda=\left(\lambda_{1}, \ldots, \lambda_{n}\right)$.
We define the following $n \times n$ matrices:

$$
\begin{gathered}
A=\left(\lambda_{1} b, \lambda_{2} b, \ldots, \lambda_{n} b\right) \\
B=(0, \ldots, 0, b, 0, \ldots, 0) \\
C=\left(\lambda_{1} d, \lambda_{2} d, \ldots, \lambda_{n} d\right) \\
D=(0, \ldots, 0, d, 0, \ldots, 0)
\end{gathered}
$$

where 0 denotes a column of n zeros.
For each $i=1, \ldots, n$, we have that $b_{i}$ satisfies the following conditions

$$
\begin{gathered}
t b_{i}=-t d_{i-1}+(1+t) d_{i}-d_{i+1}, \quad i \neq p \\
t b_{p}=-t d_{p-1}+(1+t) d_{p}-d_{p+1}+t \\
\sum_{i=1}^{n} \lambda_{i} b_{i}=-\left(1+d_{p}+t\right)
\end{gathered}
$$

setting any undefined $d_{j}$ equal zero.
For any choice $\lambda=\left(\lambda_{1}, \ldots, \lambda_{n}\right)$, we get a linear representation

$$
\psi_{\lambda}: A\left(E_{n+1, p}\right) \rightarrow G l_{2 n}(R)
$$

where R is the field of rational fractions in $\mathrm{n}+1$ indeterminates $\mathbb{Q}\left(t, d_{1}, \ldots, d_{n}\right)$.

$$
\psi_{\lambda}\left(\sigma_{i}\right) \rightarrow\left(\begin{array}{ll}
I_{n} & 0 \\
R_{i} & J_{i}
\end{array}\right)
$$

$$
\psi_{\lambda}(\delta) \rightarrow\left(\begin{array}{cc}
I_{n}+A & B \\
C & I_{n}+D
\end{array}\right)
$$

For more details, see (T.E.Brendle, 2002).

## 4. Reducibility of $\psi_{\lambda}: A\left(E_{3,1}\right) \rightarrow \boldsymbol{G} L_{4}(\mathbb{C})$

Having defined Perron's representation, we set $n=2$ and $p=1$ to get the following vectors. $b=\binom{b_{1}}{b_{2}}, d=\binom{d_{1}}{d_{2}}$, and $\lambda=\left(\lambda_{1}, \lambda_{2}\right)$.
We get the following $2 \times 2$ matrices $\quad A=\left(\begin{array}{ll}\lambda_{1} b_{1} & \lambda_{2} b_{1} \\ \lambda_{1} b_{2} & \lambda_{2} b_{2}\end{array}\right), \quad B=\left(\begin{array}{ll}b_{1} & 0 \\ b_{2} & 0\end{array}\right), \quad C=\left(\begin{array}{ll}\lambda_{1} d_{1} & \lambda_{2} d_{1} \\ \lambda_{1} d_{2} & \lambda_{2} d_{2}\end{array}\right), \quad$ and $\quad D=$ $\left(\begin{array}{ll}d_{1} & 0 \\ d_{2} & 0\end{array}\right)$.

Simple computations show that the parameters satisfy the following equations:

- $t b_{2}=-t d_{1}+(1+t) d_{2}$
- $t b_{1}=(1+t) d_{1}-d_{2}+t$
- $\lambda_{1} b_{1}+\lambda_{2} b_{2}=-\left(1+t+d_{1}\right)$

Having defined the $2 \times 2$ matrices $A, B, C$ and $D$, we obtain the multiparameter representation $A\left(E_{3,1}\right)$. This representation is of degree 4 . We specialize the parameters $\lambda_{1}, \lambda_{2}, b_{1}, b_{2}, d_{1}, d_{2}, t$ to values in $\mathbb{C}-\{0\}$. We further assume that $t \neq-1$ and $d_{2}=-t$. The representation $\psi_{\lambda}: A\left(E_{3,1}\right) \rightarrow G L_{4}(\mathbb{C})$ is defined as follows:

$$
\begin{aligned}
& \psi_{\lambda}\left(\sigma_{1}\right)=\left(\begin{array}{cccc}
1 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 \\
t & 0 & -t & 1 \\
0 & 0 & 0 & 1
\end{array}\right), \\
& \psi_{\lambda}\left(\sigma_{2}\right)=\left(\begin{array}{llll}
1 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & t & t & -t
\end{array}\right),
\end{aligned}
$$

and

$$
\psi_{\lambda}(\delta)=\left(\begin{array}{cccc}
1+\lambda_{1} b_{1} & \lambda_{2} b_{1} & b_{1} & 0 \\
\lambda_{1} b_{2} & \lambda_{2} b_{2}+1 & b_{2} & 0 \\
\lambda_{1} d_{1} & \lambda_{2} d_{1} & 1+d_{1} & 0 \\
-t \lambda_{1} & -t \lambda_{2} & -t & 1
\end{array}\right) .
$$

The graph $E_{3,1}$ has 3 vertices $\sigma_{1}, \sigma_{2}$ and $\delta$. Since $p=1$, it follows that the vertex $\delta$ has a common edge with $\sigma_{p}=\sigma_{1}$. Therefore, the following relations are satisfied.

$$
\begin{align*}
& \sigma_{1} \sigma_{2} \sigma_{1}=\sigma_{2} \sigma_{1} \sigma_{2}  \tag{1}\\
& \sigma_{2} \delta=\delta \sigma_{2}  \tag{2}\\
& \sigma_{1} \delta \sigma_{1}=\delta \sigma_{1} \delta \tag{3}
\end{align*}
$$

We note that relation (1) is actually Artin's braid relation of the classical braid group, $B_{3}$ having $\sigma_{1}$ and $\sigma_{2}$ as standard generators. This assures that a representation of $A\left(E_{3,1}\right)$ yields a representation of $B_{3}$.
Lemma 1 The representation $\psi_{\lambda}: A\left(E_{3,1}\right) \rightarrow G L_{4}(\mathbb{C})$ is reducible.
Proof. For simplicity, we write $\sigma_{i}$ instead of $\psi_{\lambda}(k)$, where $k$ is a generator of $A\left(E_{3,1}\right)$. The subspace $S=$ $\left\langle e_{1}+\frac{b_{2}}{b_{1}} e_{2}, e_{3}, e_{4}\right\rangle$ is an invariant subspace of dimension 3. To see this:

1. $\sigma_{1}\left(e_{1}+\frac{b_{2}}{b_{1}} e_{2}\right)=e_{1}+\frac{b_{2}}{b_{1}} e_{2}+t e_{3} \in S$
2. $\sigma_{2}\left(e_{1}+\frac{b_{2}}{b_{1}} e_{2}\right)=e_{1}+\frac{b_{2}}{b_{1}} e_{2}+t \frac{b_{2}}{b_{1}} e_{3} \in S$
3. $\delta\left(e_{1}+\frac{b_{2}}{b_{1}} e_{2}\right)=\left(1+\lambda_{1} b_{1}+\lambda_{2} b_{2}\right) e_{1}+\left(\lambda_{1} b_{2}+\frac{b_{2}}{b_{1}}\left(\lambda_{2} b_{2}+1\right)\right) e_{2}+$

$$
\begin{aligned}
& \left(\lambda_{1} d_{1}+\frac{b_{2}}{b_{1}} \lambda_{2} d_{1}\right) e_{3}+\left(-t \lambda_{1}+\frac{-t b_{2}}{b_{1}} \lambda_{2}\right) e_{4} \\
= & \left(1+\lambda_{1} b_{1}+\lambda_{2} b_{2}\right)\left(e_{1}+\frac{b_{2}}{b_{1}} e_{2}\right)+\left(\lambda_{1} d_{1}+\frac{b_{2}}{b_{1}} \lambda_{2} d_{1}\right) e_{3}+ \\
& \left(-t \lambda_{1}+\frac{-t b_{2}}{b_{1}} \lambda_{2}\right) e_{4} \in S
\end{aligned}
$$

4. $\sigma_{1} e_{3}=-t e_{3} \in S$
5. $\sigma_{2} e_{3}=e_{3}+t e_{4} \in S$
6. $\delta e_{3}=b_{1}\left(e_{1}+\frac{b_{2}}{b_{1}} e_{2}\right)+\left(1+d_{1}\right) e_{3}-t e_{4} \in S$
7. $\sigma_{1} e_{4}=e_{3}+e_{4} \in S$
8. $\sigma_{2} e_{4}=-t e_{4} \in S$
9. $\delta e_{4}=e_{4} \in S$

## 5. On the Irreducibility of $\psi_{\lambda}^{\prime}: A\left(E_{3,1}\right) \rightarrow G L_{3}(\mathbb{C})$

We consider the representation $\psi_{\lambda}: A\left(E_{3,1}\right) \rightarrow G L_{4}(\mathbb{C})$ restricted to the basis $e_{1}, e_{1}+\frac{b_{2}}{b_{1}} e_{2}, e_{3}$, and $e_{4}$. The matrix of $\sigma_{1}$ becomes

$$
\psi_{\lambda}\left(\sigma_{1}\right)=\left(\begin{array}{cccc}
1 & 0 & t & 0 \\
0 & 1 & t & 0 \\
0 & 0 & -t & 0 \\
0 & 0 & 1 & 1
\end{array}\right)
$$

Similarly, we determine the matrices of $\sigma_{2}$ and $\delta$. It is easy to see that the first column of the matrices of all generators is $(1,0,0,0,0)^{T}$, where $T$ is the transpose. We thus reduce our representation to a 3-dimensional one by deleting the first row and the first column to get $\psi_{\lambda}^{\prime}: A\left(E_{3,1}\right) \rightarrow G L_{3}(\mathbb{C})$. The representation is defined as follows:

$$
\begin{aligned}
& \psi_{\lambda}^{\prime}\left(\sigma_{1}\right)=\left(\begin{array}{ccc}
1 & t & 0 \\
0 & -t & 0 \\
0 & 1 & 1
\end{array}\right), \\
& \psi_{\lambda}^{\prime}\left(\sigma_{2}\right)=\left(\begin{array}{ccc}
1 & 0 & \frac{t b_{2}}{b_{1}} \\
0 & 1 & t \\
0 & 0 & -t
\end{array}\right),
\end{aligned}
$$

and

$$
\psi_{\lambda}^{\prime}(\delta)=\left(\begin{array}{ccc}
1+\lambda_{1} b_{1}+\lambda_{2} b_{2} & \lambda_{1} d_{1}+\frac{b_{2}}{b_{1}} \lambda_{2} d_{1} & -t \lambda_{1}+\frac{-t b_{2}}{b_{1}} \lambda_{2} \\
b_{1} & 1+d_{1} & -t \\
0 & 0 & 1
\end{array}\right)
$$

We then diagonalize the matrix corresponding to $\psi_{\lambda}^{\prime}\left(\sigma_{1}\right)$ by an invertible matrix, say $T$, and conjugate the matrices of $\psi_{\lambda}^{\prime}\left(\sigma_{2}\right)$ and $\psi_{\lambda}^{\prime}(\delta)$ by the same matrix $T$. The invertible matrix $T$ is given by

$$
T=\left(\begin{array}{ccc}
0 & 1 & t \\
0 & 0 & -1-t \\
1 & 0 & 1
\end{array}\right)
$$

In fact, a computation shows that

$$
T^{-1} \psi_{\lambda}^{\prime}\left(\sigma_{1}\right) T=\left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & 1 & 0 \\
0 & 0 & -t
\end{array}\right)
$$

After conjugation, we get

$$
T^{-1} \psi_{\lambda}^{\prime}\left(\sigma_{2}\right) T=\left(\begin{array}{ccc}
\frac{-t^{2}}{1+t} & 0 & \frac{-\left(1+t+t^{2}\right)}{1+t} \\
\frac{t\left(b_{2}+b_{1} t+b_{2} t\right)}{b_{1}(1+t)} & 1 & \frac{t\left(b_{2}+b_{1} t+b_{2} t\right)}{b_{1}(1+t)} \\
\frac{-t}{1+t} & 0 & \frac{1}{1+t}
\end{array}\right)
$$

$T^{-1} \psi_{\lambda}^{\prime}(\delta) T=$
$\left(\begin{array}{ccc}\frac{1}{1+t} & \frac{b_{1}}{1+t} & \frac{t}{1+t} \\ -t\left(\lambda_{1}+\frac{b_{2} \lambda_{2}}{b_{1}}+\frac{t}{1+t}\right) & 1+\lambda_{2} b_{2}+b_{1}\left(\lambda_{1}+\frac{t}{1+t}\right) & \frac{\left(-t+b_{1} t-d_{1}(1+t)\right)\left[b_{2} \lambda_{2}(1+t)+b_{1}\left(\lambda_{1}+t+\lambda_{1} t\right)\right]}{b_{1}(1+t)} \\ \frac{t}{1+t} & \frac{b_{1}}{1+t} & \frac{1}{1+t}\end{array}\right)$.
The entries of the matrices $T^{-1} \psi_{\lambda}^{\prime}\left(\sigma_{2}\right) T$ and $T^{-1} \psi_{\lambda}^{\prime}(\delta) T$ are well-defined since we assume in our work that $t \neq-1$. For simplicity, we denote $T^{-1} \psi_{\lambda}^{\prime}\left(\sigma_{1}\right) T$ by $\psi_{\lambda}^{\prime}\left(\sigma_{1}\right), T^{-1} \psi_{\lambda}^{\prime}\left(\sigma_{2}\right) T$ by $\psi_{\lambda}^{\prime}\left(\sigma_{2}\right)$, and $T^{-1} \psi_{\lambda}^{\prime}(\delta) T$ by $\psi_{\lambda}^{\prime}(\delta)$.
We now prove some propositions to determine a sufficient condition for irreducibility of $\psi_{\lambda}^{\prime}: A\left(E_{3,1}\right) \rightarrow G L_{3}(\mathbb{C})$.
Proposition $2 t\left(b_{2}+b_{1} t+b_{2} t\right)+\left(1+t+t^{2}\right)\left(\lambda_{2} b_{2}(1+t)+b_{1}\left(\lambda_{1}+t+\lambda_{1} t\right)\right)=-(t+1)^{2}\left(t^{2}+1\right)$
Proof. The proof easily follows by considering the following relations:

- $t b_{2}=-t d_{1}-t(1+t)$
- $t b_{1}=(1+t) d_{1}+2 t$
- $\lambda_{1} b_{1}+\lambda_{2} b_{2}=-\left(1+t+d_{1}\right)$

Proposition 3 The two expressions $1+t+t^{2}$ and $b_{1} t+b_{2} t+b_{2}$ cannot be both equal to zeros.
Proof. We assume, for contradiction, that both are equal to zeros.
By substituting $t b_{2}=-t d_{1}-t(1+t)$ and $t b_{1}=(1+t) d_{1}+2 t \quad$ in $b_{1} t+b_{2} t+b_{2}=0$, we get $-t\left(1+t+t^{2}\right)=-t^{2}$. By assuming that $1+t+t^{2}=0$, we get that $t=0$, a contradiction.
Proposition $4-t+b_{1} t-d_{1}(1+t) \neq 0$.
Proof. Assume, for contradiction, that $-t+b_{1} t-d_{1}(1+t)=0$. Having that $b_{1} t=(1+t) d_{1}+2 t$, we get $-t+b_{1} t-d_{1}(1+t)=-t+(1+t) d_{1}+t+t-d_{1}(1+t)=t$. This implies that $t=0$, a contradiction.
We use Proposition 2, Proposition 3 and Proposition 4 to prove the following Lemma. We recall that all the indeterminates used in defining the representations are specialized to non zero complex numbers and, in addition, the complex number associated with $t$ is not equal to -1 .

Lemma 5 If $t \neq \pm i$, then any non zero subspace $S$, which is invariant under the action of the representation $\psi_{\lambda}^{\prime}: A\left(E_{3,1}\right) \rightarrow G l_{3}(\mathbb{C})$ containing the standard unit vector $e_{3}$, must be the whole space $\mathbb{C}^{3}$.
Proof. We have that $\psi_{\lambda}^{\prime}\left(\sigma_{2}\right)\left(e_{3}\right)=\frac{-\left(1+t+t^{2}\right)}{1+t} e_{1}+\frac{t\left(b_{2}+b_{1} t+b_{2} t\right)}{b_{1}(1+t)} e_{2}+\frac{1}{1+t} e_{3} \in S$.
Since $e_{3} \in S$, it follows that $\frac{-\left(1+t+t^{2}\right)}{1+t} e_{1}+\frac{t\left(b_{2}+b_{1} t+b_{2} t\right)}{b_{1}(1+t)} e_{2} \in S$.
Moreover,
$\psi_{\lambda}^{\prime}(\delta)\left(e_{3}\right)=$
$\frac{-t+b_{1} t-d_{1}(1+t)}{1+t} e_{1}+\frac{\left(-t+b_{1} t-d_{1}(1+t)\right)\left[\lambda_{2} b_{2}(1+t)+b_{1}\left(\lambda_{1}+t+\lambda_{1} t\right)\right]}{b_{1}(1+t)} e_{2}+\frac{1+t+d_{1}(1+t)+t-b_{1} t}{1+t} e_{3} \in S$.
This also implies that

$$
\begin{equation*}
\frac{-t+b_{1} t-d_{1}(1+t)}{1+t} e_{1}+\frac{\left(-t+b_{1} t-d_{1}(1+t)\right)\left[\lambda_{2} b_{2}(1+t)+b_{1}\left(\lambda_{1}+t+\lambda_{1} t\right)\right]}{b_{1}(1+t)} e_{2} \in S \tag{2}
\end{equation*}
$$

Having proved that $1+t+t^{2}$ and $b_{1} t+b_{2} t+b_{2}$ can't both be zeros, we consider the following cases:
Case 1. $1+t+t^{2}=0$
By Proposition 3 and (1), we get that $e_{2} \in S$. By Proposition 4 and (2), we get that $e_{1} \in S$. Thus, S is the whole space.
Case 2. $1+t+t^{2} \neq 0$
Let us multiply (1) by $-t+b_{1} t-d_{1}(1+t)$ which is proved not to be zero in Proposition 4 . We also multiply (2) by $1+t+t^{2} \neq 0$. If we add the obtained equations, we get

$$
\begin{equation*}
\frac{-t+b_{1} t-d_{1}(1+t)}{b_{1}(1+t)}\left[t\left(b_{2}+b_{1} t+b_{2} t\right)+\left(1+t+t^{2}\right)\left(\lambda_{2} b_{2}(1+t)+b_{1}\left(\lambda_{1}+t+\lambda_{1} t\right)\right)\right] e_{2} \in S \tag{3}
\end{equation*}
$$

By Proposition 2, we have that $t\left(b_{2}+b_{1} t+b_{2} t\right)+\left(1+t+t^{2}\right)\left(\lambda_{2} b_{2}(1+t)+b_{1}\left(\lambda_{1}+t+\lambda_{1} t\right)\right)=-(t+1)^{2}\left(t^{2}+1\right)$. Assuming that $t \neq-1$ and $t \neq \pm i$, we get $\left[t\left(b_{2}+b_{1} t+b_{2} t\right)+\left(1+t+t^{2}\right)\left(\lambda_{2} b_{2}(1+t)+b_{1}\left(\lambda_{1}+t+\lambda_{1} t\right)\right) \neq 0\right.$.
By Proposition 4 and by (3), we get

$$
e_{2} \in S
$$

From (1) we conclude that

$$
e_{1} \in S
$$

Thus, $S$ is the whole space $\mathbb{C}^{3}$.
Next, we present the following theorem which gives a sufficient condition for irreducibility of $\psi_{\lambda}^{\prime}: A\left(E_{3,1}\right) \rightarrow$ $G L_{3}(\mathbb{C})$.
Theorem 6 If $t \neq \pm i$, then the representation $\psi_{\lambda}^{\prime}: A\left(E_{3,1}\right) \rightarrow G L_{3}(\mathbb{C})$ is irreducible.
Proof. Let $S$ be a non zero proper subspace of $\mathbb{C}^{3}$, which is invariant under the action of $\psi_{\lambda}^{\prime}$. By Lemma 5 , we have that $e_{3} \notin S$.
Then $S$ is one of the following subspaces:

- $S=\left\langle e_{1}\right\rangle$
- $S=\left\langle e_{2}\right\rangle$
- $S=\left\langle e_{1}+u e_{2}\right\rangle$, where $u \in \mathbb{C}^{*}$
- $S=\left\langle e_{1}, e_{2}\right\rangle$

Case 1. $S=\left\langle e_{1}\right\rangle$. We have that $\psi_{\lambda}^{\prime}\left(\sigma_{2}\right)\left(e_{1}\right) \in S$. This implies that $\left(\begin{array}{c}\frac{-t^{2}}{1+t} \\ \frac{t\left(b_{2}+b_{1} t+b_{2} t\right)}{b_{1}(1+t)} \\ \frac{-t}{1+t}\end{array}\right) \in S$. This gives a contradiction because $t \neq 0$.

Case 2. $S=\left\langle e_{2}\right\rangle$. We have that $\psi_{\lambda}^{\prime}(\delta)\left(e_{2}\right) \in S$. This implies that $\left(\begin{array}{c}\frac{b_{1}}{1+t} \\ 1+\lambda_{2} b_{2}+b_{1}\left(\lambda_{1}+\frac{t}{1+t}\right) \\ \frac{b_{1}}{1+t}\end{array}\right) \in S$. This gives a contradiction since $b_{1} \neq 0$.

Case 3. $S=\left\langle e_{1}+u e_{2}\right\rangle, u \in \mathbb{C}^{*}$. We have that $\psi_{\lambda}^{\prime}\left(\sigma_{2}\right)\left(e_{1}+u e_{2}\right) \in S$. This implies that $\left(\begin{array}{c}\frac{-t^{2}}{1+t} \\ \frac{t\left(b_{2}+b_{1} t+b_{2} t\right)}{b_{1}(1+t)}+u \\ \frac{-t}{1+t}\end{array}\right) \in S$. This gives a contradiction since $t \neq 0$.
Case 4. $S=\left\langle e_{1}, e_{2}\right\rangle$.
We have that $\psi_{\lambda}^{\prime}\left(\sigma_{2}\right)\left(e_{1}\right) \in S$. This implies that $\left(\begin{array}{c}\frac{-t^{2}}{1+t} \\ \frac{t\left(b_{2}+b_{1} t+b_{2} t\right)}{b_{1}(1+t)} \\ \frac{-t}{1+t}\end{array}\right) \in S$. This gives a contradiction since $t \neq 0$.
Therefore, we conclude that the representation is irreducible because there is no proper non zero invariant subspace under the action of $\psi_{\lambda}^{\prime}$.
We now give a necessary condition for irreducibility.
Theorem 7 If $t= \pm i$, then the subspace $\left\langle e_{1}, e_{3}\right\rangle$ is a proper invariant subspace.
Proof.

1. $\psi_{\lambda}^{\prime}\left(\sigma_{1}\right)\left(e_{1}\right)=e_{1} \in S$.
2. $\psi_{\lambda}^{\prime}\left(\sigma_{2}\right)\left(e_{1}\right)=\left(\begin{array}{c}\frac{-t^{2}}{1+t} \\ \frac{t\left(b_{2}+b_{1} t+b_{2} t\right)}{b_{1}(1+t)} \\ \frac{-t}{1+t}\end{array}\right)=a e_{1}+b e_{3}$, where $a$ and $b \in \mathbb{C}-\{0\}$.

Here, we have $a=\frac{-t^{2}}{1+t}, b=\frac{-t}{1+t}$, and $\frac{t\left(b_{2}+b_{1} t+b_{2} t\right)}{b_{1}(1+t)}=0$. This is true since $t^{2}=-1$.
3. $\psi_{\lambda}^{\prime}(\delta)\left(e_{1}\right)=\left(\begin{array}{c}\frac{1}{1+t} \\ -t\left(\lambda_{1}+\frac{b_{2} \lambda_{2}}{b_{1}}+\frac{t}{1+t}\right. \\ \frac{t}{1+t}\end{array}\right)=a e_{1}+b e_{3}$,
where $a$ and $b$ are given by $a=\frac{1}{1+t}$ and $b=\frac{t}{1+t}$.
4. $\psi_{\lambda}^{\prime}\left(\sigma_{1}\right)\left(e_{3}\right)=-t e_{3} \in S$.
5. $\psi_{\lambda}^{\prime}\left(\sigma_{2}\right)\left(e_{3}\right)=\left(\begin{array}{c}\frac{-\left(1+t+t^{2}\right.}{1+t} \\ \frac{t\left(b_{2}+b_{1} t+b_{2} t\right)}{b_{1}(1+t)} \\ \frac{1}{1+t}\end{array}\right)=a e_{1}+b e_{3}$,
where $a=\frac{-\left(1+t+t^{2}\right)}{1+t}$ and $b=\frac{1}{1+t}$.
6. $\psi_{\lambda}^{\prime}(\delta) e_{3}=\left(\begin{array}{c}\frac{t}{1+t} \\ \frac{\left(-t+b_{1} t-d_{1}(1+t)\right)\left(b_{2} \lambda_{2}(1+t)+b_{1}\left(\lambda_{1}+t+\lambda_{1} t\right)\right)}{b_{1}(1+t)} \\ \frac{1}{1+t}\end{array}\right)=a e_{1}+b e_{3}$,
where $a$ and $b$ are given by $a=\frac{t}{1+t}$ and $b=\frac{1}{1+t}$.
By Proposition 2, we have $\frac{\left(-t+b_{1} t-d_{1}(1+t)\right)\left(b_{2} \lambda_{2}(1+t)+b_{1}\left(\lambda_{1}+t+\lambda_{1} t\right)\right)}{b_{1}(1+t)}=0$.
Thus, we have determined a necessary and sufficient condition for irreducibility.
Theorem 8 Let $\lambda_{1}, \lambda_{2}, b_{1}, b_{2}, d_{1}, t \in \mathbb{C}-\{0\}$ and $t \neq-1$. The representation $\psi_{\lambda}^{\prime}: A\left(E_{3,1}\right) \rightarrow G L_{3}(\mathbb{C})$ is irreducible if and only if $t \neq \pm i$.

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